

HIGHER-ORDER CORRECTIONS AT THE LHC: CURRENT STATUS AND PROSPECTS

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HOCTools-II

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- 1 Introduction
- 2 Feynman diagrams
- 3 The NLO revolution
- 4 Towards higher precision: NNLO $2 \rightarrow 3$, N³LO $2 \rightarrow 2$
- 5 Summary - Discussion

Higgs

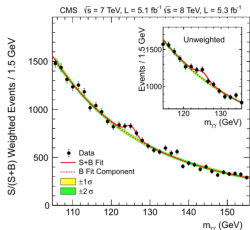


Fig. 3. The diphoton invariant mass distribution with each event weighted by the $S/(S+B)$ value of its category. The lines represent the fitted background and signal, and the coloured bands represent the ± 1 and ± 2 standard deviation uncertainties in the background estimate. The inset shows the central part of the unweighted invariant mass distribution. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this Letter.)

Gravitational wave

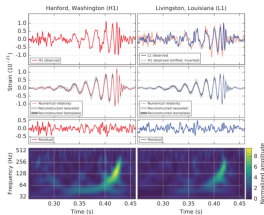


Fig. 1. The gravitational-wave event GW150914 observed by the LIGO Hanford (H1, left column panels) and Livingston (L2, right column panels) detectors. These are shown relative to September 14, 2015 at 09:50:43 UTC. For visualization, all time series are filtered with a 35–700 Hz bandpass filter to suppress large fluctuations outside the detector’s most sensitive frequency band, and bandpassed filters to remove the rising instrumental noise lines seen in the Fig. 1 spectra. Top row, left: 800 strain. Top row, right: L1 strain. GW150914 arrived first at L1 and 6.7 ms later at H1. For a visual comparison, the H1 data are also shown, offset in time by this amount and inverted in sign for the detector’s relative orientation. Second row: Gravitational-wave strain projected onto each detector on the 35–700 Hz band. Solid lines show a numerical relativity waveform for a system with parameters consistent with those measured from GW150914 [12–15] combined with 90% of an independent calibration band at [12]. Shaded areas show 90% confidence regions for two independent waveform reconstructions. One (dark grey) models the signal using binary black hole amplitude variations [26]. The other (light grey) does not use an amplitude model, but instead calculates the strain signal as a linear combination of one-Gaussian wavelets [10,14]. These two reconstructions have a 0.8% overlap, as shown in [10]. Third row: Residuals after subtracting the fitted numerical relativity waveform from the filtered detector strain. Bottom row: A spectrogram representation [12] of the strain data, showing the signal frequency increasing over time.

BH Horizon

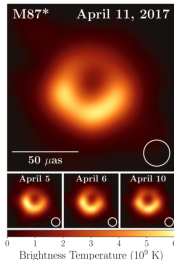
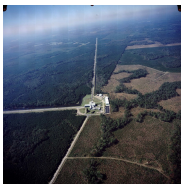


Figure 3. Top: EHT image of M87* from observations on 2017 April 11 as a representative example of the images collected in the 2017 campaign. The image is the average of three different imaging methods after convolving each with a circular Gaussian kernel to match resolution. The largest of the three kernels (29 μas FWHM) is shown in the lower right. The image is shown in units of brightness temperature, $T_b = S^2/2k\lambda^2$, where S is the flux density, λ is the observing wavelength, k_B is the Boltzmann constant, and Ω is the solid angle of the resolution element. Bottom: similar images taken over different days showing the stability of the basic image structure and the significance among different days. North is up and east is to the left.

LHC



LIGO



EHT

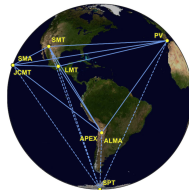
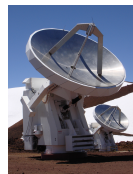


Figure 1. Eight stations of the EHT 2017 campaign over six geographic locations as viewed from the equatorial plane. Solid lines represent mutual visibility on MST ($n=12$ directions). The dashed lines were used for the calibration source JC279 (see Pages III and IV).



LHC

LIGO

EHT

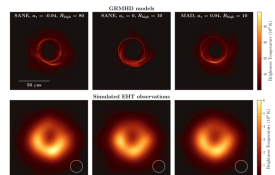
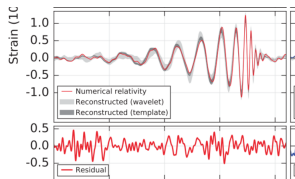
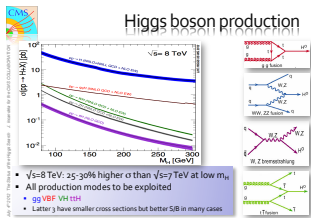
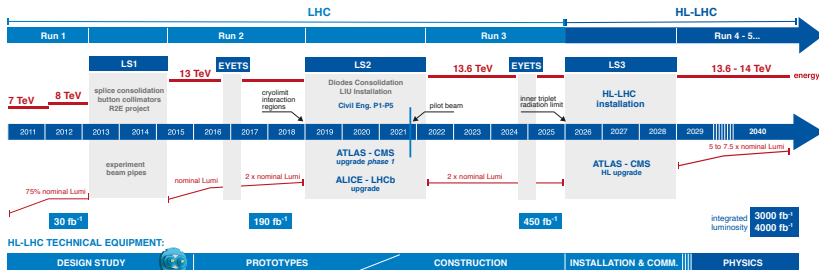


Figure 4. Top: Three example results of some of the best fitting waveforms from the large library of GRMHD simulations for April 11 corresponding to different spin parameters and accretion rates. Bottom: The same numerical models, processed through VLBI correlation pipelines with the same pipeline, showing characteristic and subtle parameters as the April 11 one and imaged in the same way as Figure 3. Note that although the fit to the observations is equally good in the three cases, they were in radically different physical scenarios. This highlights that a single good fit does not imply that a model is preferred over others (see Figure 1).

Faint signals; Patience; Theory



LHC / HL-LHC Plan



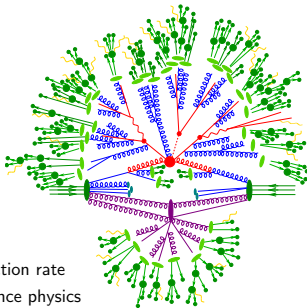
HL-LHC CIVIL ENGINEERING:

DEFINITION	EXCAVATION	BUILDINGS
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Factorization

Collins, Soper, Sterman '85-'89

- ▶ Calculate
 - ▶ Scattering probability
 - ▶ Gluon emission probability
- ▶ Measure
 - ▶ Long distance interactions
 - ▶ Particle decay rates



Divide et Impera

- ▶ Quantity of interest: Total interaction rate
- ▶ Convolution of short & long distance physics

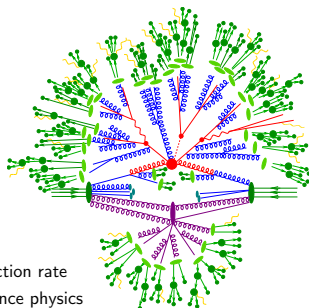
$$\sigma_{p_1, p_2 \rightarrow X} = \sum_{i, j \in \{q, g\}} \int dx_1 dx_2 \underbrace{f_{p_1, i}(x_1, \mu_F^2) f_{p_2, j}(x_2, \mu_F^2)}_{\text{long distance physics}} \underbrace{\hat{\sigma}_{ij \rightarrow X}(x_1 x_2, \mu_F^2)}_{\text{short distance physics}}$$

QCD as a perturbative quantum field theory

Factorization

Collins, Soper, Stermann '85-'89

- ▶ Calculate
 - ▶ Scattering probability
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Divide et Impera

- ▶ Quantity of interest: Total interaction rate
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$$\sigma_{p_1 p_2 \rightarrow X} = \sum_{i,j \in \{q,g\}} \int dx_1 dx_2 \underbrace{f_{p_1,i}(x_1, \mu_F^2) f_{p_2,j}(x_2, \mu_F^2)}_{\text{long distance physics}} \underbrace{\hat{\sigma}_{ij \rightarrow X}(x_1 x_2, \mu_F^2)}_{\text{short distance physics}}$$

QCD as a perturbative quantum field theory

Lattice QCD results:

→ C. Alexandrou, et al. Phys. Rev. Lett. **121**, 112001 (2018) [arXiv:1803.02685 [hep-lat]].

Leading Order

How to avoid Feynman diagrams

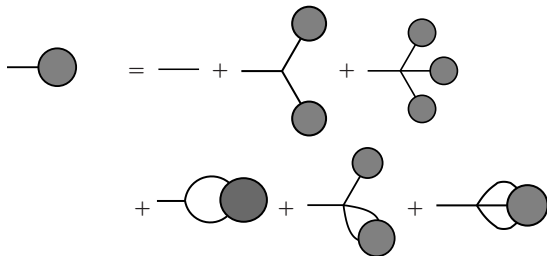
MadGraph

→ T. Stelzer and W. F. Long, *Comput. Phys. Commun.* **81**, 357 (1994)

- 1999 HELAC: The first code to calculate recursively tree-order amplitudes for (practically) arbitrary number of particles

From Feynman Diagrams to recursive equations: taming the $n!$

- 1999 HELAC: The first code to calculate recursively tree-order amplitudes for (practically) arbitrary number of particles



Unfortunately not so much on the second line !

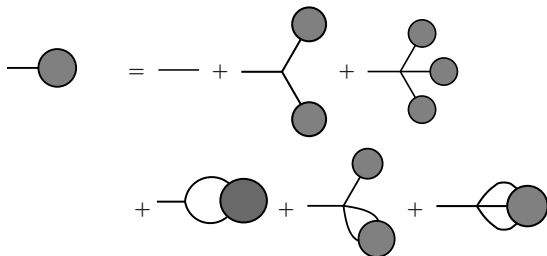
From Feynman Diagrams to recursive equations: taming the $n!$

- 1999 HELAC: The first code to calculate recursively tree-order amplitudes for (practically) arbitrary number of particles

→ A. Kanaki and C. G. Papadopoulos, *Comput. Phys. Commun.* **132** (2000) 306 [arXiv:hep-ph/0002082].

→ F. A. Berends and W. T. Giele, *Nucl. Phys. B* **306** (1988) 759.

→ F. Caravaglios and M. Moretti, *Phys. Lett. B* **358** (1995) 332.



Unfortunately not so much on the second line !

From Feynman graphs ...

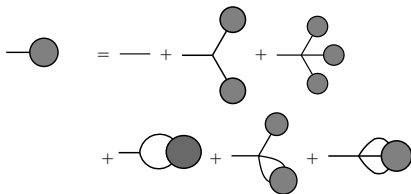
$gg \rightarrow ng$	2	3	4	5	6	7	8	9
# FG	4	25	220	2,485	34,300	559,405	10,525,900	224,449,225

TAMING THE BEAST ...

From Feynman graphs ...

$gg \rightarrow ng$	2	3	4	5	6	7	8	9
# FG	4	25	220	2,485	34,300	559,405	10,525,900	224,449,225

to Dyson-Schwinger recursion! Helac-Phegas



$gg \rightarrow ng$	2	3	4	5	6	7	8	9
#	5	15	35	70	126	210	330	495

NLO

Don't make integrals, make integrands !

What do we need for an NLO calculation ?

$$p_1, p_2 \rightarrow p_3, \dots, p_{m+2}$$

$$\begin{aligned} \sigma_{NLO} &= \int_m d\Phi_m |M_m^{(0)}|^2 J_m(\Phi) \quad \leftarrow LO \\ &+ \int_m d\Phi_m 2\text{Re}(M_m^{(0)*} M_m^{(1)}(\epsilon_{UV}, \epsilon_{IR})) J_m(\Phi) \quad \leftarrow Virtual \\ &+ \int_{m+1} d\Phi_{m+1} |M_{m+1}^{(0)}|^2 J_{m+1}(\Phi) \quad \leftarrow Real \end{aligned}$$

$J_m(\Phi)$ jet function: Infrared safeness $J_{m+1} \rightarrow J_m$

What do we need for an NLO calculation ?

$$p_1, p_2 \rightarrow p_3, \dots, p_{m+2}$$

$$\begin{aligned} \sigma_{NLO} &= \int_m d\Phi_m^{D=4} (|M_m^{(0)}|^2 + 2\text{Re}(M_m^{(0)*} M_m^{(CT)}(\epsilon_{UV}))) J_m(\Phi) \\ &+ \int_m d\Phi_m^{D=4} 2\text{Re}(M_m^{(0)*} M_m^{(1)}(\epsilon_{UV}, \epsilon_{IR})) J_m(\Phi) \\ &+ \int_{m+1} d\Phi_{m+1}^{D=4-2\epsilon_{IR}} |M_{m+1}^{(0)}|^2 J_{m+1}(\Phi) \end{aligned}$$

IR and UV divergencies, Four-Dimensional-Helicity scheme; scale dependence μ_R

What do we need for an NLO calculation ?

$$p_1, p_2 \rightarrow p_3, \dots, p_{m+2}$$

$$\begin{aligned} \sigma_{NLO} &= \int_m d\Phi_m J_m(\Phi) \\ &+ \int_m d\Phi_m 2\text{Re}(M_m^{(0)*} M_m^{(1)}(\epsilon_{UV}, \epsilon_{IR})) J_m(\Phi) \\ &+ \int_{m+1} d\Phi_{m+1} |M_{m+1}^{(0)}|^2 J_{m+1}(\Phi) \end{aligned}$$

QCD factorization— μ_F Collinear counter-terms when PDF are involved

THE ONE LOOP PARADIGM

basis of scalar integrals:

known already before NLO-R; remember this is not the case for higher orders

→ G. 't Hooft and M. J. G. Veltman, Nucl. Phys. B **153** (1979) 365.

→ Z. Bern, L. J. Dixon and D. A. Kosower, Nucl. Phys. B **412** (1994) 751

→ G. Passarino and M. J. G. Veltman, Nucl. Phys. B **160** (1979) 151.

→ Z. Bern, L. J. Dixon, D. C. Dunbar and D. A. Kosower, Nucl. Phys. B **425** (1994) 217.

$$\mathcal{A} = \sum d_{i_1 i_2 i_3 i_4} \text{[square diagram]} + \sum c_{i_1 i_2 i_3} \text{[triangle diagram]} + \sum b_{i_1 i_2} \text{[bubble diagram]} + \sum a_{i_1} \text{[self-energy diagram]} + R$$

$a, b, c, d \rightarrow$ cut-constructible part

$R \rightarrow$ rational terms

$$\mathcal{A} = \sum_{l \in \{0, 1, \dots, m-1\}} \int \frac{\mu^{(4-d)d^d q}}{(2\pi)^d} \frac{\bar{N}_l(\bar{q})}{\prod_{i \in l} \bar{D}_i(\bar{q})}$$

THE OLD “MASTER” FORMULA

$$\begin{aligned}
 \mathcal{A} \rightarrow \int \frac{N(q)}{\bar{D}_0 \bar{D}_1 \cdots \bar{D}_{m-1}} &= \sum_{i_0 < i_1 < i_2 < i_3}^{m-1} d(i_0 i_1 i_2 i_3) \int \frac{1}{\bar{D}_{i_0} \bar{D}_{i_1} \bar{D}_{i_2} \bar{D}_{i_3}} \\
 &+ \sum_{i_0 < i_1 < i_2}^{m-1} c(i_0 i_1 i_2) \int \frac{1}{\bar{D}_{i_0} \bar{D}_{i_1} \bar{D}_{i_2}} \\
 &+ \sum_{i_0 < i_1}^{m-1} b(i_0 i_1) \int \frac{1}{\bar{D}_{i_0} \bar{D}_{i_1}} \\
 &+ \sum_{i_0}^{m-1} a(i_0) \int \frac{1}{\bar{D}_{i_0}} \\
 &+ \text{rational terms}
 \end{aligned}$$


OPP “MASTER” FORMULA - I

General expression for the 4-dim $N(q)$ at the integrand level in terms of D_i

$$\begin{aligned} N(q) &= \sum_{i_0 < i_1 < i_2 < i_3}^{m-1} \left[d(i_0 i_1 i_2 i_3) + \tilde{d}(q; i_0 i_1 i_2 i_3) \right] \prod_{i \neq i_0, i_1, i_2, i_3}^{m-1} D_i \\ &+ \sum_{i_0 < i_1 < i_2}^{m-1} \left[c(i_0 i_1 i_2) + \tilde{c}(q; i_0 i_1 i_2) \right] \prod_{i \neq i_0, i_1, i_2}^{m-1} D_i \\ &+ \sum_{i_0 < i_1}^{m-1} \left[b(i_0 i_1) + \tilde{b}(q; i_0 i_1) \right] \prod_{i \neq i_0, i_1}^{m-1} D_i \\ &+ \sum_{i_0}^{m-1} \left[a(i_0) + \tilde{a}(q; i_0) \right] \prod_{i \neq i_0}^{m-1} D_i \end{aligned}$$

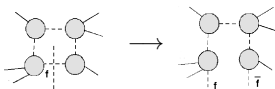
THE ONE-LOOP CALCULATION IN A NUTSHELL

The computation of $pp(p\bar{p}) \rightarrow e^+ \nu_e \mu^- \bar{\nu}_\mu b \bar{b}$ involves up to six-point functions. The most generic integrand has therefore the form

$$\mathcal{A}(q) = \sum \underbrace{\frac{N_i^{(6)}(q)}{\bar{D}_{i_0} \bar{D}_{i_1} \cdots \bar{D}_{i_5}}}_{\text{hexagon}} + \underbrace{\frac{N_i^{(5)}(q)}{\bar{D}_{i_0} \bar{D}_{i_1} \cdots \bar{D}_{i_4}}}_{\text{pentagon}} + \underbrace{\frac{N_i^{(4)}(q)}{\bar{D}_{i_0} \bar{D}_{i_1} \cdots \bar{D}_{i_3}}}_{\text{square}} + \underbrace{\frac{N_i^{(3)}(q)}{\bar{D}_{i_0} \bar{D}_{i_1} \bar{D}_{i_2}}}_{\text{triangle}} + \dots$$


In order to apply the OPP reduction, HELAC evaluates numerically the numerators $N_i^6(q), N_i^5(q), \dots$ with the values of the loop momentum q provided by CutTools

- generates all inequivalent partitions of 6,5,4,3... blobs attached to the loop, and check all possible flavours (and colours) that can be consistently running inside
- hard-cuts the loop (q is fixed) to get a $n+2$ tree-like process



The R_2 contributions (rational terms) are calculated in the same way as the tree-order amplitude, taking into account *extra vertices*

→ BlackHat, MadGraph, RECOLA, OpenLoops

THE ONE-LOOP CALCULATION IN A NUTSHELL

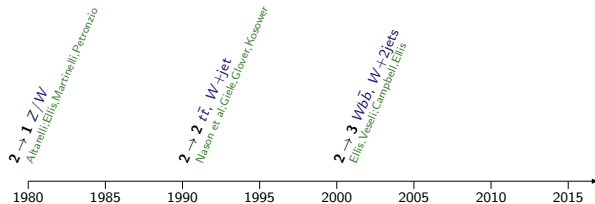
Institute of Nuclear Physics "Demokritos", Bergische Universität Wuppertal, Institute of Nuclear Physics PAN, RWTH Aachen University

	Content
<h2>HELAC-NLO & Associated Tools</h2>	
Projects	
HELAC-PHEGAS - A generator for all parton level processes in the Standard Model	
HELAC-DIPOLES - Dipole formalism for the arbitrary helicity eigenstates of the external partons	
HELAC-ILoop - A program for numerical evaluation of QCD virtual corrections to scattering amplitudes	
ONELOOP - A program for the evaluation of one-loop scalar functions	
CUTTOOLS - A program implementing the OPP reduction method to compute one-loop amplitudes	
FARNI - A program for importance sampling and density estimation	
KALEU - A general-purpose parton-level phase space generator	
HELAC-ONIA - An automatic matrix element generator for heavy quarkonium physics	
-- top --	
People	
Giuseppe Bevilacqua	
Michael Czakon	
Maria Vittoria Garzelli	
Andreas van Hameren	
Adam Kardos	
Yiannis Malamou	
Costas G. Papadopoulos	
Roberto Pittau	
Malgorzata Worek	
Hua-Sheng Shao	
-- top --	
Contact us	
If you have a question, comment, suggestion or bug report, please e-mail us at:	
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erdosshen@gmail.com	
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Last modified by Malgorzata Worek Thursday, January 10th, 2013	

Proof of concept: the first NLO public code

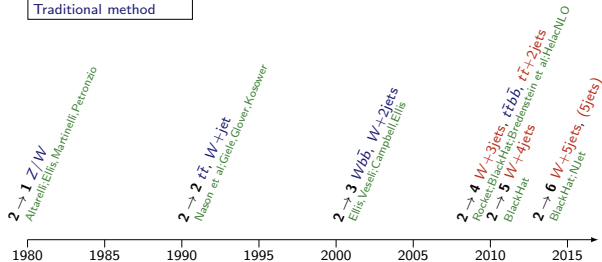
G. P. Salam, PoS ICHEP 2010, 556 (2010) [arXiv:1103.1318 [hep-ph]]

The NLO revolution



The NLO revolution

Unitarity based method
Traditional method



BlackHat → Berger, Bern, Dixon, Febres Cordero, Forde, Ita, Kosower, Maitre

HelacNLO → Bevilacqua, Czakon, Papadopoulos, Pittau, Worek

NJet → Badger, Biedermann, Uwer, Yundin

Rockett → Ellis, Melnikov, Zanderighi

MadGraph:

→ J. Alwall et al., JHEP **1407** (2014) 079 [arXiv:1405.0301 [hep-ph]].

OpenLoops:

→ F. Cascioli, P. Maierhofer and S. Pozzorini, Phys. Rev. Lett. **108**, 111601 (2012) [arXiv:1111.5206 [hep-ph]].

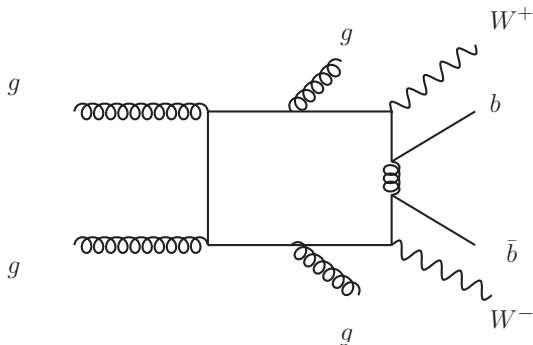
The NLO wishlist

Process ($V \in \{Z, W, \gamma\}$)	Status
1. $pp \rightarrow VV$ jet	WW jet completed by Dittmaier/Kallweit/Uwer; Campbell/Ellis/Zanderighi ZZ jet completed by Binoth/Gleisberg/Karg/Kauer/Sanguinetti WZ jet, $W\gamma$ jet completed by Campanario et al.
2. $pp \rightarrow$ Higgs+2 jets	NLO QCD to the gg channel completed by Campbell/Ellis/Zanderighi NLO QCD+EW to the VBF channel completed by Cicolini/Denner/Dittmaier Interference QCD-EW in VBF channel
3. $pp \rightarrow V V V$	ZZZ completed by Lazopoulos/Melnikov/Petriello and WWZ by Hankele/Zeppenfeld see also Binoth/Ossola/Papadopoulos/Pittau VBFNLO meanwhile also contains WWW, ZZW, ZZZ, WW γ , ZZ γ , WZ γ , $W\gamma\gamma$, $Z\gamma\gamma$, $\gamma\gamma\gamma$, $W\gamma\gamma$
4. $pp \rightarrow t\bar{t} b\bar{b}$	relevant for $t\bar{t}H$, computed by Bredenstein/Denner/Dittmaier/Pozzorini and Bevilacqua/Czakon/Papadopoulos/Pittau/Worek
5. $pp \rightarrow V+3$ jets	$W+3$ jets calculated by the Blackhat/Sherpa and Rocket collaborations
6. $pp \rightarrow t\bar{t}+2$ jets	$Z+3$ jets by Blackhat/Sherpa relevant for $t\bar{t}H$, computed by Bevilacqua/Czakon/Papadopoulos/Worek
7. $pp \rightarrow VV b\bar{b}$,	Pozzorini et al. Bevilacqua et al.
8. $pp \rightarrow VV+2$ jets	W^+W^-+2 jets, W^+W^-+2 jets, relevant for VBF $H \rightarrow VV$ VBF contributions by (Bozzi/Jäger/Oleari/Zeppenfeld
9. $pp \rightarrow b\bar{b}b\bar{b}$	Binoth et al.
10. $pp \rightarrow V+4$ jets	top pair production, various new physics signatures Blackhat/Sherpa: $W+4$ jets, $Z+4$ jets see also HEJ for $W+4$ jets
11. $pp \rightarrow Wb\bar{b}j$	top, new physics signatures, Reina/Schutzmeier
12. $pp \rightarrow t\bar{t}t\bar{t}$	various new physics signatures, Bevilacqua/Worek
$pp \rightarrow W\gamma\gamma$ jet	Campanario/Englert/Rauch/Zeppenfeld
$pp \rightarrow 4/5$ jets	Blackhat+Sherpa/NJets



- ▶ NLO calculations requested by LHC experimenters
- ▶ List constructed in 2005
- ▶ Calculations completed 2012

→ G. Bevilacqua, M. Lupattelli, D. Stremmer and M. Worek, [arXiv:2212.04722 [hep-ph]].



NLO $2 \rightarrow 6$ ($2 \rightarrow 8$ including leptonic W^\pm decays)

Towards higher precision: NNLO and beyond

→ A. Huss, J. Huston, S. Jones and M. Pellen, [arXiv:2207.02122 [hep-ph]].

process	known	desired
$pp \rightarrow H$	$N^3\text{LO}_{\text{HTL}}$	$N^4\text{LO}_{\text{HTL}}$ (incl.) $\text{NNLO}_{\text{QCD}}^{(h,e)}$
	$\text{NNLO}_{\text{QCD}}^{(l)}$	
	$N^{(1,1)}\text{LO}_{\text{QCD}\otimes\text{EW}}^{(\text{HTL})}$	
	NLO_{QCD}	
$pp \rightarrow H + j$	NNLO_{HTL}	$\text{NNLO}_{\text{HTL}} \otimes \text{NLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$
	NLO_{QCD}	
	$N^{(1,1)}\text{LO}_{\text{QCD}\otimes\text{EW}}$	
$pp \rightarrow H + 2j$	$\text{NLO}_{\text{HTL}} \otimes \text{LO}_{\text{QCD}}$	$\text{NNLO}_{\text{HTL}} \otimes \text{NLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$ $N^3\text{LO}_{\text{QCD}}^{(\text{VBF}^*)}$ $\text{NNLO}_{\text{QCD}}^{(\text{VBF}^*)}$ $\text{NLO}_{\text{EW}}^{(\text{VBF})}$
	$N^3\text{LO}_{\text{QCD}}^{(\text{VBF}^*)}$ (incl.)	
	$\text{NNLO}_{\text{QCD}}^{(\text{VBF}^*)}$	
	$\text{NLO}_{\text{EW}}^{(\text{VBF})}$	
	NLO_{HTL}	
$pp \rightarrow H + 3j$	$\text{NLO}_{\text{QCD}}^{(\text{VBF})}$	$\text{NLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$
	NLO_{QCD}	
$pp \rightarrow VH$	$\text{NNLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$	$\text{NNLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$
	$\text{NLO}_{gg \rightarrow HZ}^{(d)}$	
$pp \rightarrow VH + j$	NNLO_{QCD}	$\text{NNLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$
	$\text{NLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$	
$pp \rightarrow HH$	$N^3\text{LO}_{\text{HTL}} \otimes \text{NLO}_{\text{QCD}}$	NLO_{EW}
$pp \rightarrow HH + 2j$	$N^3\text{LO}_{\text{QCD}}^{(\text{VBF}^*)}$ (incl.)	NNLO_{QCD}
	$\text{NNLO}_{\text{QCD}}^{(\text{VBF}^*)}$	
	$\text{NLO}_{\text{EW}}^{(\text{VBF})}$	
$pp \rightarrow HHH$	NNLO_{HTL}	
$pp \rightarrow H + t\bar{t}$	$\text{NLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$	NNLO_{QCD}
	NNLO_{QCD} (off-diag.)	
$pp \rightarrow H + t/\bar{t}$	$\text{NLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$	NNLO_{QCD}

Table 1: Precision wish list: Higgs boson final states. $N^4\text{LO}_{\text{QCD}}^{(\text{VBF}^*)}$ means a calculation using the structure function approximation. $V = W, Z$.

→ A. Huss, J. Huston, S. Jones and M. Pellen, [arXiv:2207.02122 [hep-ph]].

process	known	desired
$pp \rightarrow 2 \text{ jets}$	NNLO _{QCD}	N ³ LO _{QCD} + NLO _{EW}
	NLO _{QCD} + NLO _{EW}	
$pp \rightarrow 3 \text{ jets}$	NNLO _{QCD} + NLO _{EW}	

Table 2: Precision wish list: jet final states.

process	known	desired
$pp \rightarrow V$	$N^5\text{LO}_{\text{QCD}}$	$N^5\text{LO}_{\text{QCD}} + N^{(1,1)}\text{LO}_{\text{QCD}\otimes\text{EW}}$ $N^2\text{LO}_{\text{EW}}$
	$N^{(1,1)}\text{LO}_{\text{QCD}\otimes\text{EW}}$	
	NLO_{EW}	
$pp \rightarrow VV'$	$\text{NNLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$ $+ \text{NLO}_{\text{QCD}}$ (gg channel)	NLO_{QCD} (gg channel, w/ massive loops) $N^{(1,1)}\text{LO}_{\text{QCD}\otimes\text{EW}}$
	$\text{NNLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$	hadronic decays
$pp \rightarrow V + j$	$\text{NLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$ (QCD component)	NNLO_{QCD}
$pp \rightarrow V + 2j$	$\text{NLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$ (EW component)	
$pp \rightarrow V + b\bar{b}$	NLO_{QCD}	$\text{NNLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$
$pp \rightarrow VV' + 1j$	$\text{NLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$	NNLO_{QCD}
$pp \rightarrow VV' + 2j$	NLO_{QCD} (QCD component)	Full $\text{NLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$
	$\text{NLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$ (EW component)	
$pp \rightarrow W^+W^+ + 2j$	Full $\text{NLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$	
$pp \rightarrow W^+W^- + 2j$	$\text{NLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$ (EW component)	
$pp \rightarrow W^+Z + 2j$	$\text{NLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$ (EW component)	
$pp \rightarrow ZZ + 2j$	Full $\text{NLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$	
$pp \rightarrow VV'V''$	NLO_{QCD}	$\text{NLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$
	NLO_{EW} (w/o decays)	
$pp \rightarrow W^\pm W^+ W^-$	$\text{NLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$	
$pp \rightarrow \gamma\gamma$	$\text{NNLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$	$N^3\text{LO}_{\text{QCD}}$
$pp \rightarrow \gamma + j$	$\text{NNLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$	$N^3\text{LO}_{\text{QCD}}$
$pp \rightarrow \gamma\gamma + j$	$\text{NNLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$	
	$+ \text{NLO}_{\text{QCD}}$ (gg channel)	
$pp \rightarrow \gamma\gamma\gamma$	NNLO_{QCD}	$\text{NNLO}_{\text{QCD}} + \text{NLO}_{\text{EW}}$

Table 3: Precision wish list: vector boson final states. $V = W, Z$ and $V', V'' = W, Z, \gamma$. Full leptonic decays are understood if not stated otherwise.

→ A. Huss, J. Huston, S. Jones and M. Pellen, [arXiv:2207.02122 [hep-ph]].

process	known	desired
$pp \rightarrow t\bar{t}$	NNLO _{QCD} + NLO _{EW} (w/o decays)	
	NLO _{QCD} + NLO _{EW} (off-shell effects)	N ³ LO _{QCD}
	NNLO _{QCD} (w/ decays)	
$pp \rightarrow t\bar{t} + j$	NLO _{QCD} (off-shell effects)	NNLO _{QCD} + NLO _{EW} (w/ decays)
	NLO _{EW} (w/o decays)	
$pp \rightarrow t\bar{t} + 2j$	NLO _{QCD} (w/o decays)	NLO _{QCD} + NLO _{EW} (w/ decays)
$pp \rightarrow t\bar{t} + V'$	NLO _{QCD} + NLO _{EW} (w/o decays)	NNLO _{QCD} + NLO _{EW} (w/ decays)
$pp \rightarrow t\bar{t} + \gamma$	NLO _{QCD} (off-shell effects)	
$pp \rightarrow t\bar{t} + Z$	NLO _{QCD} (off-shell effects)	
$pp \rightarrow t\bar{t} + W$	NLO _{QCD} + NLO _{EW} (off-shell effects)	
$pp \rightarrow t/\bar{t}$	NNLO _{QCD} *(w/ decays)	NNLO _{QCD} + NLO _{EW} (w/ decays)
	NLO _{EW} (w/o decays)	
$pp \rightarrow tZj$	NLO _{QCD} + NLO _{EW} (w/ decays)	NNLO _{QCD} + NLO _{EW} (w/o decays)
$pp \rightarrow t\bar{t}\bar{t}$	NLO _{QCD} + NLO _{EW} (off-shell effects)	
	Full NLO _{QCD} + NLO _{EW} (w/o decays)	NNLO _{QCD}

Table 4: Precision wish list: top quark final states. NNLO_{QCD} * means a calculation using the structure function approximation. $V' = W, Z, \gamma$.

The two-loop frontier: $2 \rightarrow 3$

5-POINT 2-LOOP - MASSLESS: ALL FAMILIES

→ T. Gehrmann, J. M. Henn and N. A. Lo Presti, Phys. Rev. Lett. **116** (2016) no.6, 062001 [erratum: Phys. Rev. Lett. **116** (2016) no.18, 189903]

[arXiv:1511.05409 [hep-ph]].

→ C. G. Papadopoulos, D. Tommasini and C. Wever, JHEP **04** (2016), 078 [arXiv:1511.09404 [hep-ph]].

→ D. Chicherin, T. Gehrmann, J. M. Henn, P. Wasser, Y. Zhang and S. Zoia, Phys. Rev. Lett. **123** (2019) no.4, 041603

→ D. Chicherin and V. Sotnikov, JHEP **20** (2020), 167

→ S. Abreu, J. Dormans, F. Febres Cordero, H. Ita, M. Kraus, B. Page, E. Pascual, M. S. Ruf and V. Sotnikov, "Caravel: A C++ framework for the computation of multi-loop amplitudes with numerical unitarity," Comput. Phys. Commun. **267** (2021), 108069

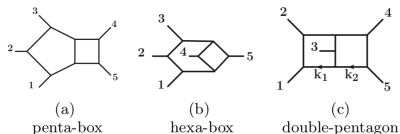


FIG. 1: Integral topologies for massless five-particle scattering at two loops.

→ J. Henn, T. Peraro, Y. Xu and Y. Zhang, "A first look at the function space for planar two-loop six-particle Feynman integrals," JHEP **03** (2022), 056

5-POINT 2-LOOP - ONE LEG OFF-SHELL: ALL FAMILIES

→ C. G. Papadopoulos, D. Tommasini and C. Wever, JHEP **04** (2016), 078 [arXiv:1511.09404 [hep-ph]].

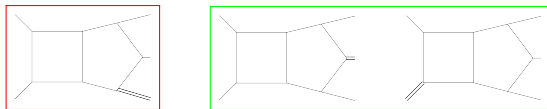
→ C. G. Papadopoulos and C. Wever, JHEP **2002** (2020) 112

→ S. Abreu, H. Ita, F. Moriello, B. Page, W. Tschernow and M. Zeng, JHEP **2011** (2020) 117

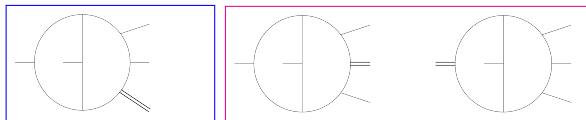
→ D. D. Canko, C. G. Papadopoulos and N. Syrrakos, JHEP **2101** (2021) 199

→ S. Abreu, H. Ita, B. Page and W. Tschernow, JHEP **03** (2022), 182 [arXiv:2107.14180 [hep-ph]].

→ A. Kardos, C. G. Papadopoulos, A. V. Smirnov, N. Syrrakos and C. Wever, [arXiv:2201.07509 [hep-ph]].



The three planar pentaboxes of the families P_1 (left), P_2 (middle) and P_3 (right) with one external massive leg.



The five non-planar families with one external massive leg.

NNLO QCD: $pp \rightarrow \gamma\gamma\gamma + X$

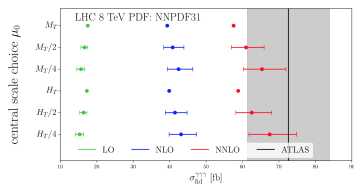


Figure 1. Predictions for the fiducial cross-section in LO (green), NLO (blue) and NNLO (red) QCD versus ATLAS data (black). Shown are predictions for six scale choices. The error bars on the theory predictions reflect scale variation only. For two of the scales only the central predictions are shown.

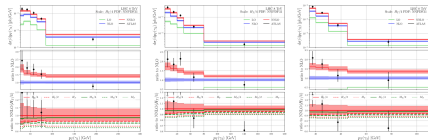


Figure 2. p_T distribution of the hardest photon γ_1 (left), γ_2 (center) and the softest one γ_3 (right). Top plot shows the absolute distribution at NNLO (red), NLO (blue) and LO (green) versus ATLAS data (black). Middle plot shows same distributions but normalized to the NLO. Bottom plot shows central NNLO predictions for 6 different scale choices (only the central scale is shown) with respect to the default choice $\mu_0 = H_T/4$. The bands represent the 7-point scale variations about the corresponding central scales.

→ H. A. Chawdhry, M. L. Czakon, A. Mitov and R. Poncelet, JHEP 2002 (2020) 057

NNLO QCD: $pp \rightarrow \gamma\gamma\gamma + X$

fiducial setup for $pp \rightarrow \gamma\gamma\gamma + X$; used in the ATLAS 8 TeV analysis of Ref. [37]		
$p_{T,\gamma_1} \geq 27 \text{ GeV}$,	$p_{T,\gamma_2} \geq 22 \text{ GeV}$,	$p_{T,\gamma_3} \geq 15 \text{ GeV}$,
$0 \leq \eta_{\gamma_1} \leq 1.37$ or $1.56 \leq \eta_{\gamma_1} \leq 2.37$,		
$\Delta R_{\gamma\gamma} \geq 0.45$,	$m_{\gamma\gamma} \geq 50 \text{ GeV}$,	Fraxione isolation with $n = 1$, $\delta_0 = 0.4$, and $E_T^{\text{eff}} = 10 \text{ GeV}$.

Table 1: Definition of phase space cuts.

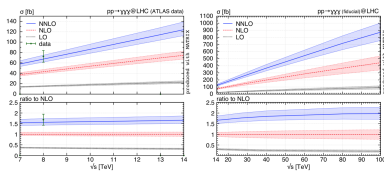


Figure 4: Fiducial cross sections for $pp \rightarrow \gamma\gamma\gamma + X$ as a function of the centre-of-mass energy at LO (black dotted), at NLO (red dashed), and at NNLO (blue, solid). The green data point at 8 TeV corresponds to the cross section measured by ATLAS in Ref. [37].

→ S. Kallweit, V. Sotnikov and M. Wiesemann, Phys. Lett. B **812** (2021) 136013

NNLO QCD: $pp \rightarrow 3\text{jets} + X$

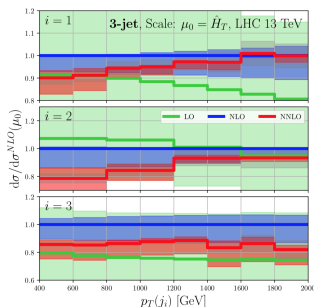
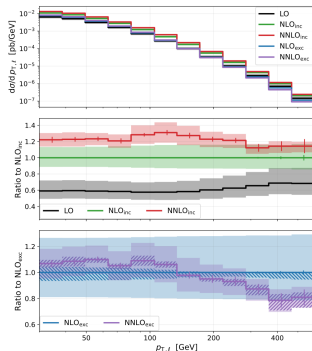


FIG. 1: The three panels show the i th leading jet transverse momentum $p_T(j_i)$ for $i = 1, 2, 3$ for the production of (at least) three jets. LO (green), NLO (blue) and NNLO (red) are shown for the central scale (solid line). 7-point scale variation is shown as a coloured band. The grey band corresponds to the uncertainty from Monte Carlo integration.

→ M. Czakon, A. Mitov and R. Poncelet, *Phys. Rev. Lett.* **127** (2021) no.15, 152001 [arXiv:2106.05331 [hep-ph]].

→ X. Chen, T. Gehrmann, N. Glover, A. Huss and M. Marcoli, [arXiv:2203.13531 [hep-ph]].

NNLO QCD: $pp \rightarrow Wb\bar{b} + X$



→ H. B. Hartanto, R. Poncelet, A. Popescu and S. Zoia, Phys. Rev. D **106** (2022) no.7, 074016 [arXiv:2205.01687 [hep-ph]].

	Comment	Complete analytic results	Public numerical code	Cross sections
$pp \rightarrow jjj$	L.c.	Abreu et al.	Abreu et al.	Chen et al., Czakon et al.
$pp \rightarrow \gamma\gamma j$	L.c.*	Agarwal et al., Chawdhry et al.	Agarwal et al.	Chawdhry et al.
$pp \rightarrow \gamma\gamma\gamma$	L.c.*	Abreu et al., Chawdhry et al.	Abreu et al.	Chawdhry et al., Kallweit et al.
$pp \rightarrow \gamma\gamma j$		Agarwal et al.		
$gg \rightarrow \gamma\gamma g$	NLO loop induced	Badger et al.	Badger et al.	Badger et al.
$pp \rightarrow Wb\bar{b}$	L.c.*, on-shell W	Badger et al.		
$pp \rightarrow W(l\nu)b\bar{b}$	L.c.	Abreu et al., Hartanto et al.		Hartanto et al.
$pp \rightarrow W(l\nu)j\bar{j}$	L.c.	Abreu et al.		
$pp \rightarrow Z(l\bar{l})j\bar{j}$	L.c.*	Abreu et al.		
$pp \rightarrow W(l\nu)\gamma j$	L.c.*	Badger et al.		
$pp \rightarrow Hb\bar{b}$	L.c., b -quark Yukawa	Badger et al.		

Table 1: Known two-loop QCD corrections for five-point scattering processes at hardon colliders. “L.c.” refers to the calculations in the leading-color approximation; “L.c.*” means that in addition non-planar L.c. contributions are omitted. All public codes employ `PentagonFunctions++` Chicherin and Sotnikov, Chicherin et al. for numerical evaluation of special functions.

The three-loop frontier: $2 \rightarrow 2$

3-LOOP CALCULATIONS

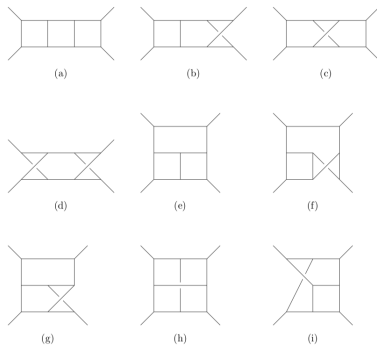


Figure 1. The nine integral families needed to describe all master integrals for three-loop massless four-particle scattering. The external legs are associated with the momenta p_1 , p_3 , p_4 and p_2 in clockwise order starting with the top left corner.

→ J. M. Henn, A. V. Smirnov and V. A. Smirnov, *JHEP* **07** (2013), 128

→ J. Henn, B. Mistlberger, V. A. Smirnov and P. Wasser, *JHEP* **04** (2020), 167

3-LOOP CALCULATIONS

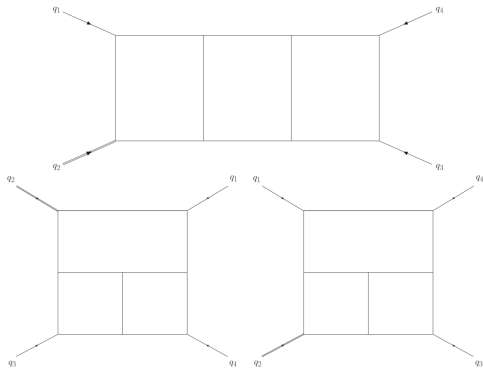


Figure 1. The F1 (top), F2 (bottom left) and F3 (bottom right) top-sector diagrams. The double line represents the massive particle and all external momenta are taken to be incoming.

→ S. Di Vita, P. Mastrolia, U. Schubert and V. Yundin, *JHEP* **09** (2014), 148

→ D. D. Canko and N. Syrrakos, *JHEP* **04** (2022), 134

3-LOOP CALCULATIONS

→ F. Caola, A. Von Manteuffel and L. Tancredi, "Diphoton Amplitudes in Three-Loop Quantum Chromodynamics," Phys. Rev. Lett. **126** (2021) no.11,

112004

→ F. Caola, A. Chakraborty, G. Gambuti, A. von Manteuffel and L. Tancredi, "Three-loop helicity amplitudes for four-quark scattering in massless QCD,"

JHEP **10** (2021), 206

→ P. Bargiela, F. Caola, A. von Manteuffel and L. Tancredi, "Three-loop helicity amplitudes for diphoton production in gluon fusion," JHEP **02** (2022), 153

→ F. Caola, A. Chakraborty, G. Gambuti, A. von Manteuffel and L. Tancredi, "Three-Loop Gluon Scattering in QCD and the Gluon Regge Trajectory,"

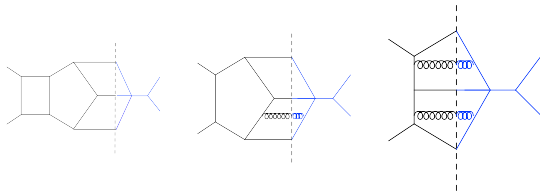
Phys. Rev. Lett. **128** (2022) no.21, 212001

→ F. Caola, A. Chakraborty, G. Gambuti, A. von Manteuffel and L. Tancredi, "Three-loop helicity amplitudes for quark-gluon scattering in QCD,"

[arXiv:2207.03503 [hep-ph]].

What do we need for an NNLO calculation ?

$$p_1, p_2 \rightarrow p_3, \dots, p_{m+2}$$



What do we need for an NNLO calculation ?

$$\begin{aligned}
 \sigma_{NNLO} &\rightarrow \int_m d\Phi_m \left(2\text{Re}(M_m^{(0)*} M_m^{(2)}) + \left| M_m^{(1)} \right|^2 \right) J_m(\Phi) && \text{VV} \\
 &+ \int_{m+1} d\Phi_{m+1} \left(2\text{Re} \left(M_{m+1}^{(0)*} M_{m+1}^{(1)} \right) \right) J_{m+1}(\Phi) && \text{RV} \\
 &+ \int_{m+2} d\Phi_{m+2} \left| M_{m+2}^{(0)} \right|^2 J_{m+2}(\Phi) && \text{RR}
 \end{aligned}$$

RV + RR → antenna-S, colorfull-NNLO, sector-improved residue subtraction, nested soft-collinear, local analytic sector subtraction, projection to born, q_T , N-jetiness

→ A. Gehrmann-De Ridder, T. Gehrmann and M. Ritzmann, JHEP **1210** (2012) 047

→ P. Bolzoni, G. Somogyi and Z. Trocsanyi, JHEP **1101** (2011) 059

→ M. Czakon and D. Heymes, Nucl. Phys. B **890** (2014) 152

→ S. Catani and M. Grazzini, Phys. Rev. Lett. **98** (2007) 222002

→ R. Boughezal, C. Focke, X. Liu and F. Petriello, Phys. Rev. Lett. **115** (2015) no.6, 062002

→ M. Cacciari, F. A. Dreyer, A. Karlberg, G. P. Salam and G. Zanderighi, Phys. Rev. Lett. **115**, no. 8, 082002 (2015)

→ F. Caola, K. Melnikov and R. Rötsch, Eur. Phys. J. C **77**, no. 4, 248 (2017)

→ L. Magnea, E. Maina, G. Pelliccioli, C. Signorile-Signorile, P. Torrielli and S. Uccirati, arXiv:1806.09570 [hep-ph].

Amplitude reduction

- Write the "OPP-type" equation at two loops

$$\frac{N(l_1, l_2; \{p_i\})}{D_1 D_2 \dots D_n} = \sum_{m=1}^{\min(n,8)} \sum_{S_{m;n}} \frac{\Delta_{i_1 i_2 \dots i_m}(l_1, l_2; \{p_i\})}{D_{i_1} D_{i_2} \dots D_{i_m}}$$

cut equations : $D_{i_1} = D_{i_2} = \dots = D_{i_m} = 0$

$\Delta_{i_1 i_2 \dots i_m}(l_1, l_2; \{p_i\}) \rightarrow$ *spurious* \oplus *ISP* – *irreducible integrals*

- Write the "OPP-type" equation at two loops

$$\frac{N(l_1, l_2; \{p_i\})}{D_1 D_2 \dots D_n} = \sum_{m=1}^{\min(n,8)} \sum_{S_{m;n}} \frac{\Delta_{i_1 i_2 \dots i_m}(l_1, l_2; \{p_i\})}{D_{i_1} D_{i_2} \dots D_{i_m}}$$

ISP-irreducible integrals \rightarrow use **IBPI** to Master Integrals

Libraries in the future: QCD2LOOP, TwOLOop

\rightarrow P. Mastrolia, T. Peraro and A. Primo, arXiv:1605.03157 [hep-ph].

\rightarrow J. Gluza, K. Kajda and D. A. Kosower, Phys. Rev. D **83** (2011) 045012

\rightarrow H. Ita, arXiv:1510.05626 [hep-th].

\rightarrow C. G. Papadopoulos, R. H. P. Kleiss and I. Malamos, PoS Corfu **2012** (2013) 019.

\rightarrow S. Abreu, J. Dormans, F. Febres Cordero, H. Ita, M. Kraus, B. Page, E. Pascual, M. S. Ruf and V. Sotnikov, Comput. Phys. Commun. **267** (2021),

108069

HELAC-2LOOP FOR AMPLITUDE CONSTRUCTION: THE ALGORITHM

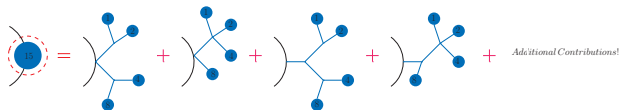
→ G. Bevilacqua, D. D. Canko, A. Kardos and C. G. Papadopoulos, J. Phys. Conf. Ser. **2105** (2021) no.5, 012010

n -particle, 2-loop Amplitude \longrightarrow $(n+2)$ -particle, 1-loop Amplitude


- 1) Definition of the flavor of the $n + 1$ and $n + 2$ particles.
- 2) Generation of the $n + 2$ color-states ($(n + 2)!$, Color-Flow Representation).
- 3) Generation of Blob-Topologies.
- 4) Cut of the topologies in the k_3 -line (middle-line) \rightarrow the 2 extra particles.
- 5) Flavor-Color Dressing of the 1-loop loop-particles.
- 6) Second cut of the blob-topology \rightarrow tree-level graph ($n + 4$ color-states).
- 7) Creation of currents contributing to the configuration (Dyson-Schwinger to blobs).
- 8) Reduction of the $n + 4$ color-states to n and identification of N_C power.
- 9) Storing of the numerator information to the Skeleton.

TWO-LOOP BLOB-TOPOLOGIES

- Binary representation for the particles: e.g. for $n = 4$, $\{1, 2, 3, 4\} \rightarrow \{1, 2, 4, 8\}$
- What a blob and its level are?



- List representation for the 3 grand blob-topologies:

Theta-topologies:  $\equiv \{\{k_1\}, \{k_2\}, \{k_3\}, \{A\}, \{B\}\}$

Infinity-topologies:  $\equiv \{\{k_1\}, \{k_2\}\}$

Dumbbell-topologies:  $\equiv \{\{k_1\}, \{k_2\}, \{C\}, \{A\}, \{B\}\}$

NUMERATORS AND NUMERICS IN 4 DIMENSIONS

Process	loop-flavors	Color	Size	Time	Numerators
$gg \rightarrow q\bar{q}$	$\{g, g_h, \bar{g}_h, q, \bar{q}\}$	full	16.1 MB	3m 14.509s	13856
$gg \rightarrow gg$	$\{g, g_h, \bar{g}_h\}$	leading	8.9 MB	15.017s	4560
$gg \rightarrow gg$	$\{g, g_h, \bar{g}_h, q, \bar{q}\}$	full	110.6 MB	6m 54.574s	89392
$gg \rightarrow ggg$	$\{g, g_h, \bar{g}_h\}$	leading	300.0 MB	21m 42.609s	81480

Comments on the skeletons:

- 1 n increase \rightarrow complexity increase
- 2 leading color to full color \rightarrow complexity increase
- 3 Timings a bit large \rightarrow Skeleton constructed only once per process!
- 4 Much numerators (some are identical) \rightarrow Room for improving efficiency!

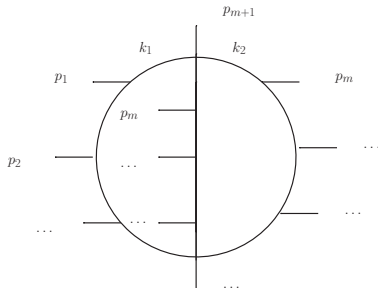
Some numerical results for numerators with gluons as external and loop particles ($h = -- \rightarrow --$)¹:

- 1 $N_{\{\{1,2\},\{12\},\{\},\{\},\{\}\}} = 17052219.315419123 + 64639250.888367772i.$
- 2 $N_{\{\{1,2\},\{4,8\},\{\},\{\},\{\}\}} = -12231870819598.090 + 5124375444085.5430i.$
- 3 $N_{\{\{1,2\},\{4\},\{8\},\{\},\{\}\}} = -1268111397619.5310 + 195312105699.88257i.$
- 4 $N_{\{\{2,1\},\{8\},\{\},\{4\},\{\}\}} = -49731029299.352333 + 15599344.440385548i.$

- Perfect agreement in cross-checks with FeynArts + FeynCalc!

Feynman Integrals

PERTURBATIVE QCD AT NNLO



$$\mathcal{N} \left(k_1, k_2; \{p_i\}_{i=1}^{m+1}, \{\varepsilon\} \right)$$

$$(k^2 - M_0^2) \left((k_1 + p_1)^2 - M_1^2 \right) \dots \left((k_1 - k_2 - p_{m+1})^2 - M_j^2 \right) \dots (k_2^2 - M_l^2)$$

THE CURRENT APPROACH

- m independent momenta, L loops, $N = L(L + 1)/2 + Lm$ scalar products
- basis composed by $D_1 \dots D_N$, allows to express all scalar products
 $D_i = (\{k_1, k_2\} + p_i)^2 - M_i^2$

- Definition

$$F[a_1, \dots, a_N] = C_L \int \frac{1}{D_1^{a_1} \dots D_N^{a_N}} \prod_{i=1}^L [d^d k_i]$$

with a_i being zero, positive or negative integers.

- Feynman parameters, Mellin-Barnes, Differential Equations
- Or numerical: SecDec, Weinzierl, Anastasiou&Sterman

→ C. Anastasiou and G. Sterman, arXiv:1812.03753 [hep-ph].

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→F. V. Tkachov, Phys. Lett. B **100** (1981) 65.

→K. G. Chetyrkin and F. V. Tkachov, Nucl. Phys. B **192** (1981) 159.

IBP identities:

$$\int d^d k d^d l \frac{\partial}{\partial \{k^\mu, l^\mu\}} \left(\frac{\{k^\mu, l^\mu, v^\mu\}}{D_1^{a_1} \dots D_N^{a_N}} \right) = 0$$

reduce *all* Feynman Integrals to a finite subset → **Master Integrals**.

$$F[a_1, \dots, a_N] = \sum_i R_i(\{p\}, d) G_i[a'_1, \dots, a'_N]$$

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→ Z. Bern, L. J. Dixon and D. A. Kosower, Phys. Lett. B **302** (1993) 299.

→ V. A. Smirnov, Phys. Lett. B **460** (1999) 397

→ T. Gehrmann and E. Remiddi, Nucl. Phys. B **580** (2000) 485 [hep-ph/9912329].

→ J. M. Henn, Phys. Rev. Lett. **110** (2013) 25, 251601 [arXiv:1304.1806 [hep-th]].

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→ S. Borowka, G. Heinrich, S. P. Jones, M. Kerner, J. Schlenk and T. Zirke, *Comput. Phys. Commun.* **196** (2015) 470

→ S. Becker, C. Reuschle and S. Weinzierl, *JHEP* **1012** (2010) 013

→ C. Anastasiou and G. Sterman, arXiv:1812.03753 [hep-ph].

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DIFFERENTIAL EQUATIONS APPROACH

- The integral is a function of external momenta, so one can set-up differential equations by differentiating and using **IBP**

$$F[a_1, \dots, a_N] \rightarrow G[a'_1, \dots, a'_N]$$

$$p_j^\mu \frac{\partial}{\partial p_i^\mu} G[a_1, \dots, a_n] \rightarrow \sum C_{b_1, \dots, b_n} F[b_1, \dots, b_n] \rightarrow \sum C_{a'_1, \dots, a'_n} G[a'_1, \dots, a'_n]$$

- Find the proper basis**; Bring the system of equations in a form suitable to express the MI in terms of GPs

$$\begin{aligned} \partial_m f(\varepsilon, \{x_i\}) &= \varepsilon A_m(\{x_i\}) f(\varepsilon, \{x_i\}) \\ \partial_m A_n - \partial_n A_m &= 0 \quad [A_m, A_n] = 0 \end{aligned}$$

★ f not MI!

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- Boundary conditions**: expansion by regions or regularity conditions.

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DIFFERENTIAL EQUATIONS APPROACH

- Iterated Integrals

→ K. T. Chen, *Iterated path integrals*, *Bull. Amer. Math. Soc.* **83** (1977) 831

- Multiple Polylogarithms, Symbol algebra

- Goncharov Polylogarithms

$$\mathcal{G}(a_n, \dots, a_1, x) = \int_0^x dt \frac{1}{t - a_n} \mathcal{G}(a_{n-1}, \dots, a_1, t)$$

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- Elliptic Integrals

→ L. Adams and S. Weinzierl, *Phys. Lett. B* **781** (2018), 270-278

→ J. Broedel, C. Duhr, F. Dulat, B. Penante and L. Tancredi, *JHEP* **01** (2019), 023

- Numerical approach [one-mass double-pentagon]

Generalised power series expansion

→ F. Moriello, *JHEP* **01** (2020), 150

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→ X. Liu and Y. Q. Ma, *Comput. Phys. Commun.* **283** (2023), 108565

DIFFERENTIAL EQUATIONS APPROACH

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→ A. B. Goncharov, M. Spradlin, C. Vergu and A. Volovich, Phys. Rev. Lett. **105** (2010) 151605.

→ C. Duhr, H. Gangl and J. R. Rhodes, JHEP **1210** (2012) 075 [arXiv:1110.0458 [math-ph]].

→ C. Bogner and F. Brown

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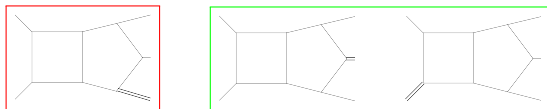
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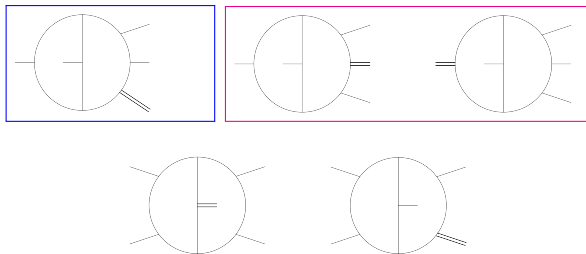
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The SDE approach

5-POINT TWO-LOOP ONE-MASS



The three planar pentaboxes of the families P_1 (left), P_2 (middle) and P_3 (right) with one external massive leg.



The five non-planar families with one external massive leg.

PENTABOX - ONE LEG OFF-SHELL: P1

→ J. M. Henn, Phys. Rev. Lett. **110** (2013) 251601

→ S. Abreu, H. Ita, F. Moriello, B. Page, W. Tschernow and M. Zeng, JHEP **2011** (2020) 117

→ D. D. Canko, C. G. Papadopoulos and N. Syrrakos, JHEP **2101** (2021) 199

$$d\vec{g} = \epsilon \sum_a d \log(W_a) \tilde{M}_a \vec{g}$$

- Also from direct differentiation of MI wrt to x . Just g in terms of FI.

$$\frac{d\vec{g}}{dx} = \epsilon \sum_b \frac{1}{x - \ell_b} M_b \vec{g}$$

- ℓ_b , are independent of x , some depending only on the reduced invariants, $\{S_{12}, S_{23}, S_{34}, S_{45}, S_{51}\}$. M_b are independent of the invariants.
- number of letters smaller than in AIMPTZ representation
- Main contribution for us from AIMPTZ: the canonical basis (+ numerics)

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$$\frac{d \log(W_a)}{dx}$$

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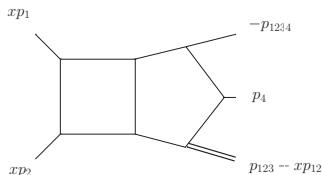
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PENTABOX - ONE LEG OFF-SHELL: P1



$$q_1 \rightarrow p_{123} - xp_{12}, \quad q_2 \rightarrow p_4, \quad q_3 \rightarrow -p_{1234}, \quad q_4 \rightarrow xp_1$$

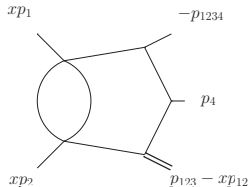
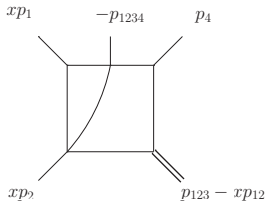
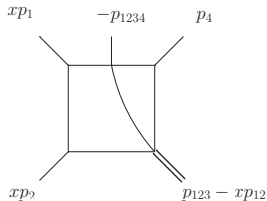
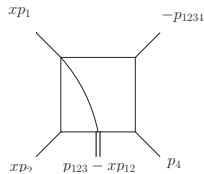
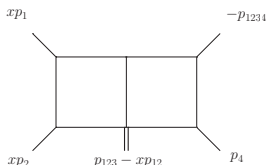
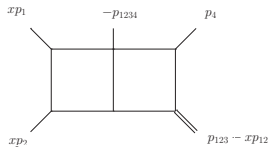
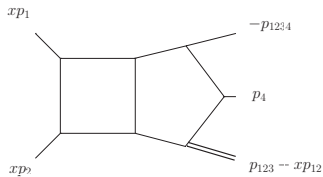
SDE parametrisation: n off-shell legs $\rightarrow n - 1$ off-shell legs + the x variable.

\rightarrow C. G. Papadopoulos, "Simplified differential equations approach for Master Integrals," JHEP **1407** (2014) 088

- p_i , $i = 1 \dots 5$, satisfy $\sum_1^5 p_i = 0$, with $p_i^2 = 0$, $i = 1 \dots 5$, $p_{i\dots j} := p_i + \dots + p_j$.
The set of independent invariants: $\{S_{12}, S_{23}, S_{34}, S_{45}, S_{51}, x\}$, with $S_{ij} := (p_i + p_j)^2$.

$$q_1^2 = (1 - x)(S_{45} - S_{12}x), \quad s_{12} = (S_{34} - S_{12}(1 - x))x, \quad s_{23} = S_{45}, \quad s_{34} = S_{51}x, \\ s_{45} = S_{12}x^2, \quad s_{15} = S_{45} + (S_{23} - S_{45})x$$

PENTABOX - ONE LEG OFF-SHELL: P1



4-POINT UP TO TWO LEGS OFF-SHELL

- J. M. Henn, K. Melnikov and V. A. Smirnov, JHEP **1405** (2014) 090
- T. Gehrmann, A. von Manteuffel, L. Tancredi and E. Weihs, JHEP **06** (2014), 032
- F. Caola, J. M. Henn, K. Melnikov and V. A. Smirnov, JHEP **1409** (2014) 043
- C. G. Papadopoulos, D. Tommasini and C. Wever, JHEP **1501** (2015) 072
- T. Gehrmann, A. von Manteuffel and L. Tancredi, JHEP **09** (2015), 128

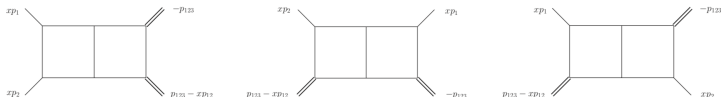


Figure 3. The parametrization of external momenta for the three planar double boxes of the families P_{12} (left), P_{13} (middle) and P_{23} (right) contributing to pair production at the LHC. All external momenta are incoming.

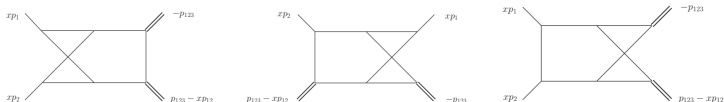


Figure 4. The parametrization of external momenta for the three non-planar double boxes of the families N_{12} (left), N_{13} (middle) and N_{34} (right) contributing to pair production at the LHC. All external momenta are incoming.

As well as planar and nonplanar double box with one off-shell leg expressed in UT basis.

PENTABOX - ONE LEG OFF-SHELL: P1-3

$$\frac{d\mathbf{g}}{dx} = \epsilon \sum_a \frac{1}{x - \ell_a} \mathbf{M}_a \mathbf{g}$$

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$$\begin{aligned} \mathbf{g} = & \epsilon^0 \mathbf{b}_0^{(0)} + \epsilon \left(\sum \mathcal{G}_a \mathbf{M}_a \mathbf{b}_0^{(0)} + \mathbf{b}_0^{(1)} \right) \\ & + \epsilon^2 \left(\sum \mathcal{G}_{ab} \mathbf{M}_a \mathbf{M}_b \mathbf{b}_0^{(0)} + \sum \mathcal{G}_a \mathbf{M}_a \mathbf{b}_0^{(1)} + \mathbf{b}_0^{(2)} \right) \\ & + \epsilon^3 \left(\sum \mathcal{G}_{abc} \mathbf{M}_a \mathbf{M}_b \mathbf{M}_c \mathbf{b}_0^{(0)} + \sum \mathcal{G}_{ab} \mathbf{M}_a \mathbf{M}_b \mathbf{b}_0^{(1)} + \sum \mathcal{G}_a \mathbf{M}_a \mathbf{b}_0^{(2)} + \mathbf{b}_0^{(3)} \right) \\ & + \epsilon^4 \left(\sum \mathcal{G}_{abcd} \mathbf{M}_a \mathbf{M}_b \mathbf{M}_c \mathbf{M}_d \mathbf{b}_0^{(0)} + \sum \mathcal{G}_{abc} \mathbf{M}_a \mathbf{M}_b \mathbf{M}_c \mathbf{b}_0^{(1)} \right. \\ & \left. + \sum \mathcal{G}_{ab} \mathbf{M}_a \mathbf{M}_b \mathbf{b}_0^{(2)} + \sum \mathcal{G}_a \mathbf{M}_a \mathbf{b}_0^{(3)} + \mathbf{b}_0^{(4)} \right) + \dots \end{aligned}$$

$$\mathcal{G}_{ab\dots} := \mathcal{G}(\ell_a, \ell_b, \dots; x)$$

PENTABOX - ONE LEG OFF-SHELL: KINEMATICAL REGIONS

- Euclidean region:

$$\left\{ S_{12} \rightarrow -2, S_{23} \rightarrow -3, S_{34} \rightarrow -5, S_{45} \rightarrow -7, S_{51} \rightarrow -11, x \rightarrow \frac{1}{4} \right\}$$

no letter l in the region $[0, x]$, all boundary terms real. [very fast GiNaC]

Family	W=1	W=2	W=3	W=4
P_1 (g_{72})	17 (14)	116 (95)	690 (551)	2740 (2066)
P_2 (g_{73})	25 (14)	170 (140)	1330 (1061)	4950 (3734)
P_3 (g_{84})	22 (12)	132 (90)	1196 (692)	4566 (2488)

TABLE: Number of GP entering in the solution. In parenthesis we give the corresponding number for the non-zero top-sector basis elements.

- with timings, running the GiNaC Interactive Shell `ginsh`, given by 1.9, 3.3, and 2 seconds for P_1 , P_2 and P_3 respectively and for a precision of 32 significant digits
- A very different canonical basis, several elements start at ϵ^4 .

SUMMARY & OUTLOOK

- **Non-planar families**

- We have completed the hexa-box families, N_1 , N_2 , N_3 .
- Checks against known results successful.
- Next task: double-pentagon families, N_4 , N_5 .

- **SDE approach: all MI up to 4-point with up to 2 off shell legs and 5-point with up to one off-shell leg.**

- **Speed-up numerical evaluation**

- Improving GPLs analytic continuation.
- Study letters ordering in physical regions, use different mappings and/or fibrations.
- Combine analytics with numerics → one-dimensional integral representation

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→ N. Syrrakos, [arXiv:2303.07395 [hep-ph]].

- **HELAC2LOOP:** generic approach to amplitude reduction and evaluation

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→ A. Kardos, C. G. Papadopoulos, A. V. Smirnov, N. Syrrakos and C. Wever, JHEP 05 (2022), 033

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→ S. Abreu, H. Ita, B. Page and W. Tschernow, JHEP **03** (2022), 182

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 - SDE@1-loop → N. Syrrakos, "One-loop Feynman integrals for 2 → 3 scattering involving many scales including internal masses," JHEP 10 (2021), 041 [arXiv:2107.02106 [hep-ph]].
 - SDE@3-loop → D. D. Canko and N. Syrrakos, "Planar three-loop master integrals for 2 → 2 processes with one external massive particle," [arXiv:2112.14275 [hep-ph]].
 - UT basis determination → more criteria as experience dictates
 - H. Frellesvig and C. G. Papadopoulos, JHEP 04 (2017), 083
 - J. Henn, B. Mistlberger, V. A. Smirnov and P. Wasser, JHEP 04 (2020), 167
 - P. Wasser, "Scattering Amplitudes and Logarithmic Differential Forms,"
 - C. Dlapa, X. Li and Y. Zhang, JHEP 07 (2021), 227
 - Boundary terms determination → for UT basis elements
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Thank you for your attention !

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